

Discerning the causality of drop breakup using massive ensembles of simulations



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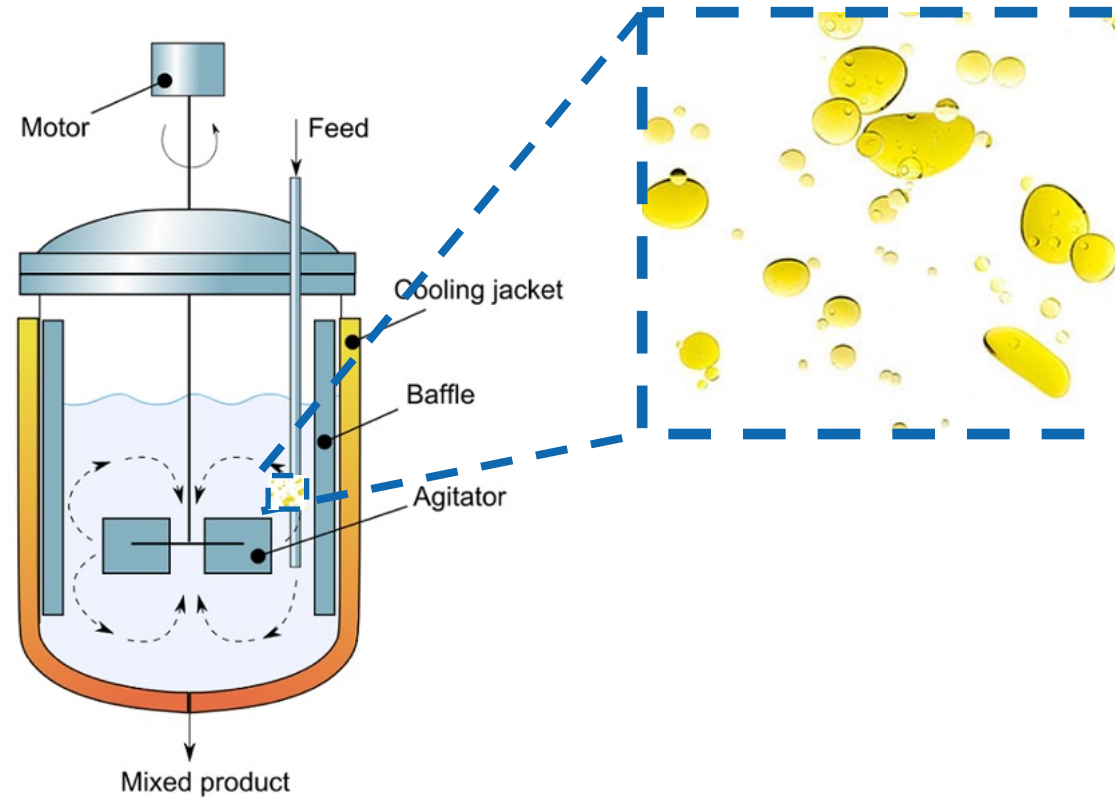
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Drop breakup

Occurs in the production of emulsions,
e.g., food, pharmaceuticals, paint

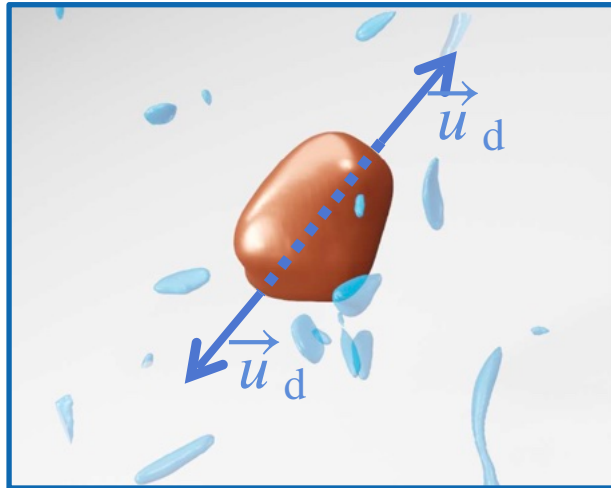


Emulsification



- Turbulent forces cause drop breakup

The Weber number



$$We = \frac{\text{kinetic energy}}{\text{surface energy}} = \frac{\frac{1}{2} \rho d^3 u_d^2}{\sigma d^2}$$

Assumptions:

- Deformation by 'eddies' of size $\sim d$,
- resistance by surface tension σ

The Kolmogorov-Hinze diameter

Kolmogorov, *Dok. Akad. Nauk.* 1949, Hinze, *AIChE J.* 1955

Kolmogorov and Hinze proposed that drops will break as long as: $We > We_c$
Rearranging gives a critical diameter

$$d_{KH} = We_c^{-3/5} \sigma^{3/5} \rho^{-3/5} \epsilon^{-2/5}, \quad We_c \approx 2$$

- ✓ Simple yes or no prediction of breakup
- ✓ Confirmed by many studies in different flows

Shinnar *J. Fluid Mech.* 1961, Vankova *et al. J. Colloid Interface Sci.* 2007, Perlekar *et al. Phys. Fluids* 2012; Cannon *et al. J. Fluid Mech.* 2023, Ni, *Annu. Rev. Fluid Mech.* 2024

- ✗ Different values of the constant We_c
- ✗ Static: No information on the dynamics, e.g. breakup time

Our questions

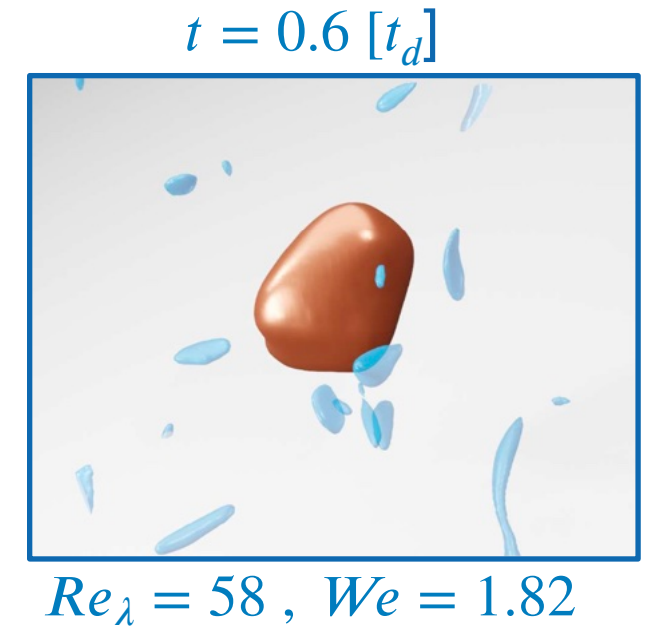
1. What is the lifetime of a drop?
2. Can we make a simple mechanistic model of drop breakup?

Method: direct numerical simulations

Vela-Martín & Avila, *J. Fluid Mech.* (2021)

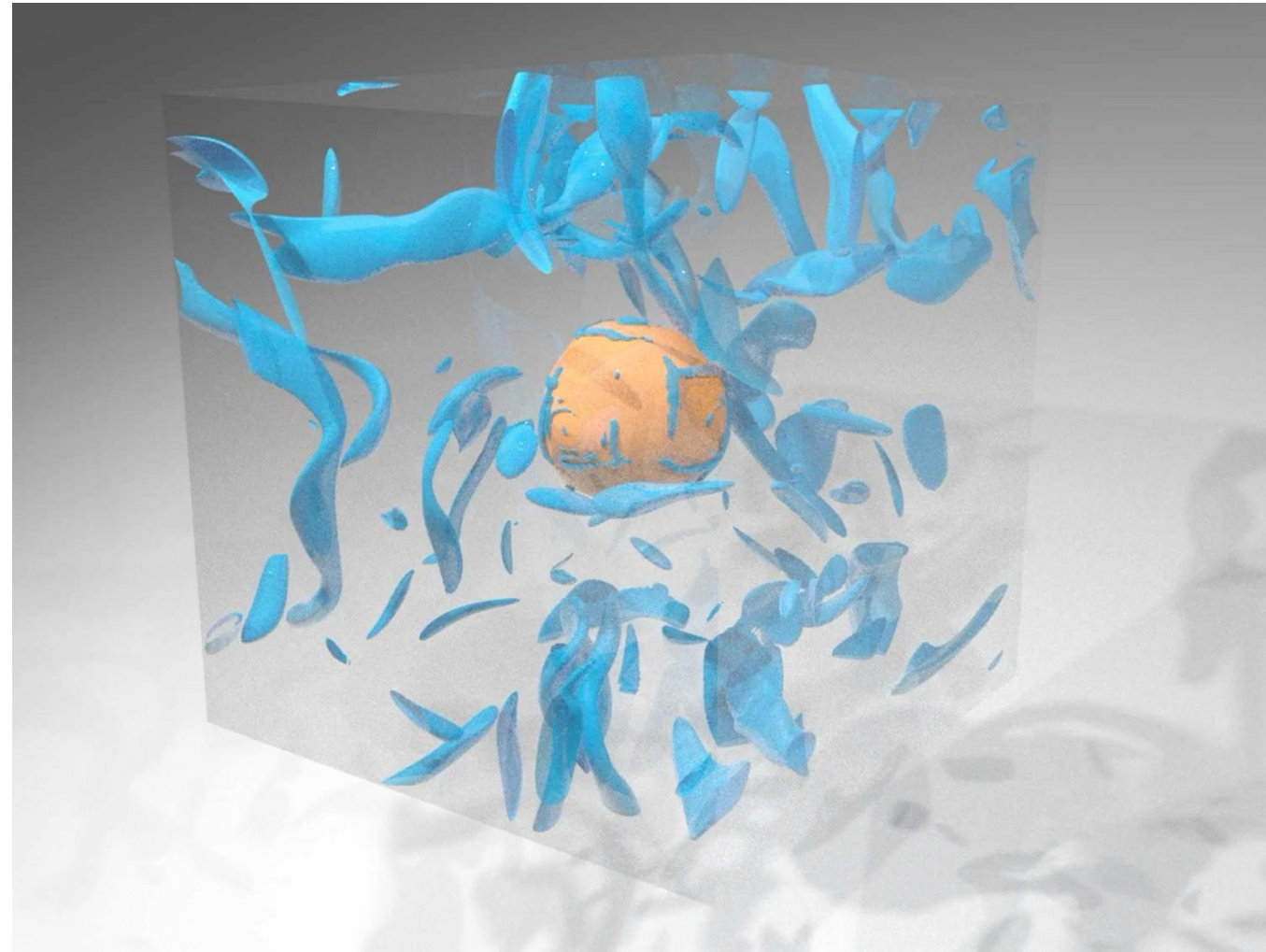
- homogeneous isotropic turbulence
- sustained by forcing on modes $k < 2$
- pseudo-spectral method
- optimised CUDA code
- 256^3 grid fits on a single GPU
- initial condition: spherical drop
- 2080 breakup simulations
- across 7 Weber numbers

	We	N
◀	1.2	180
▲	1.3	200
▼	1.4	200
▶	1.6	200
●	1.8	200
■	2.9	500
◆	3.6	600



* Jaqmin *J. Comput. Phys.* (1999), Magaletti *et al. J. Fluid Mech.* (2013)

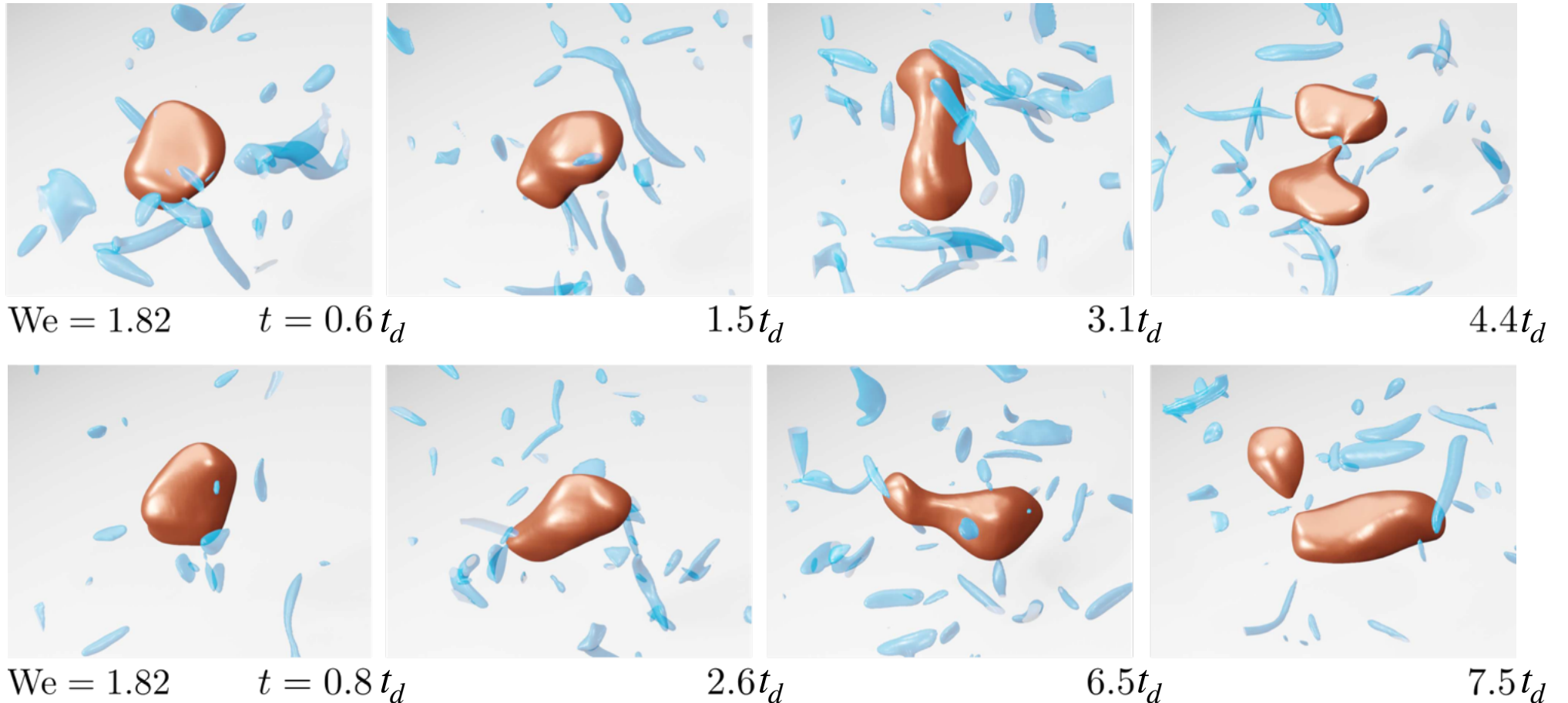
DNS of drop breakup in HIT



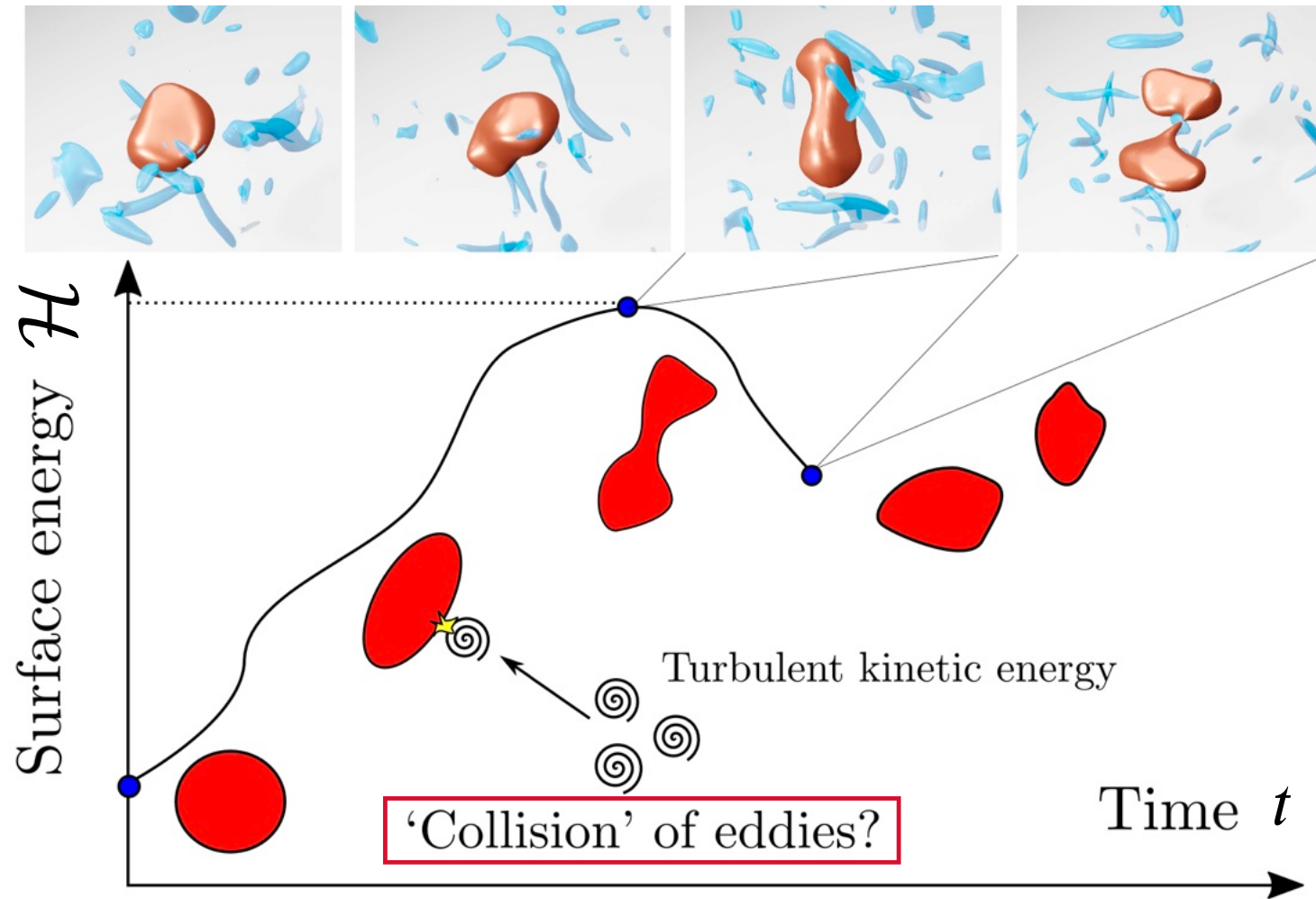
<https://youtu.be/AEsZjPGtYoA?si=gIVKadQIlz-KLPIn>

Drop breakup is a stochastic process

$$t_d = d/u_d$$



Drop deformation by 'collision' of eddies



[illegible]

Exact equations: $d_t E = \mathcal{P} - \epsilon - \langle \vartheta \rangle_S$
 $d_t \mathcal{H} = \langle \vartheta \rangle_S$

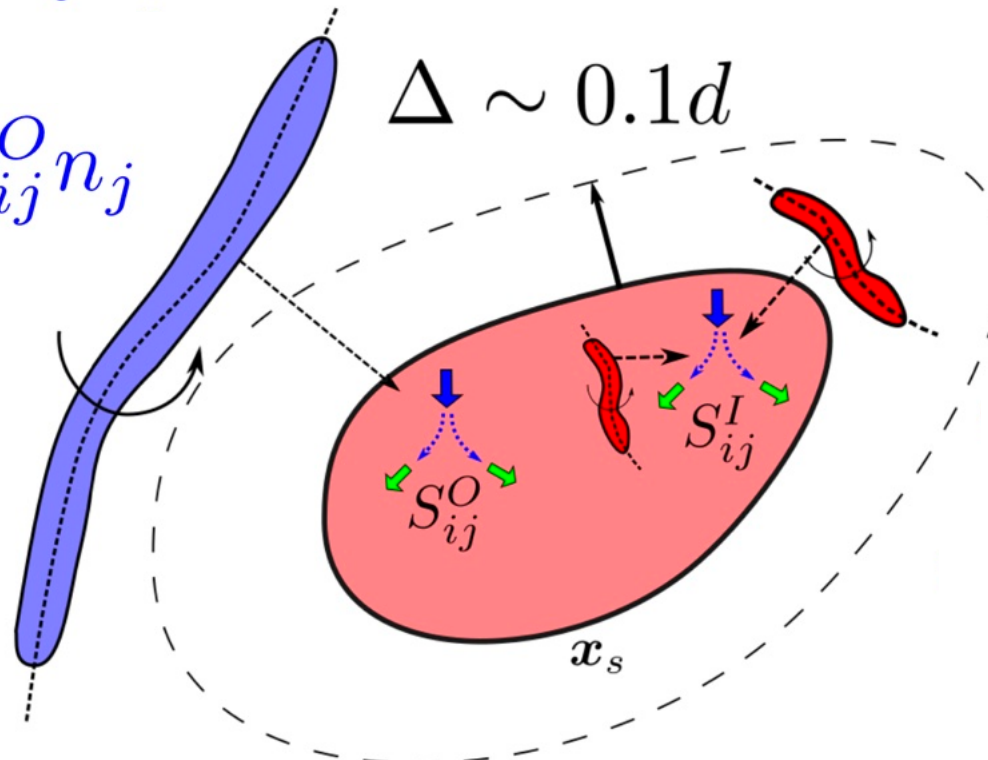
Stretching by outer/inner eddies

Decomposition of the rate-of-strain tensor (Biot-Savart law):

$$S_{ij} = S_{ij}^I + S_{ij}^O$$

$$\vartheta^O = -\sigma n_i S_{ij}^O n_j$$

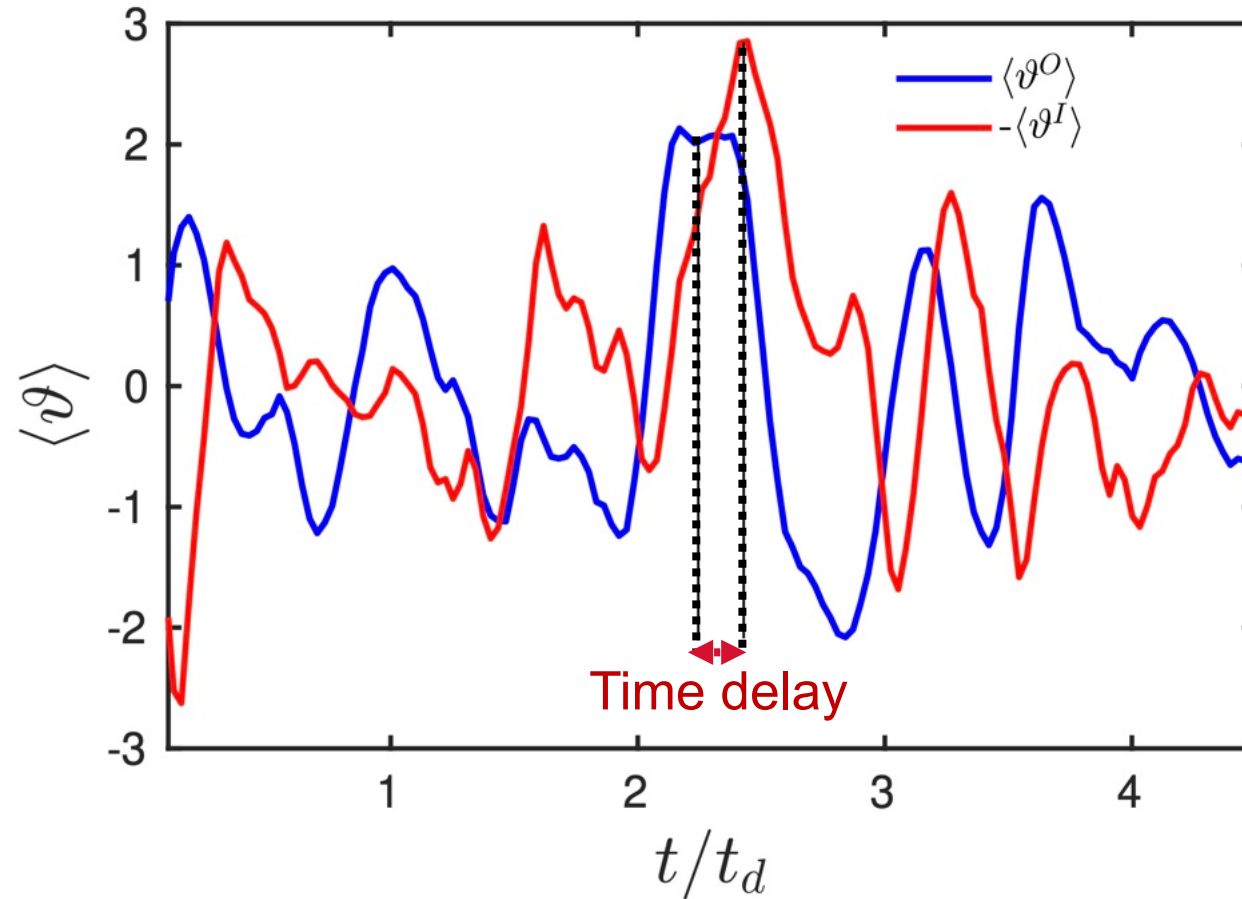
$$\vartheta^I = -\sigma n_i S_{ij}^I n_j$$



Modeling approach:
 $d_t \mathcal{H} = \vartheta_O + \vartheta_I$

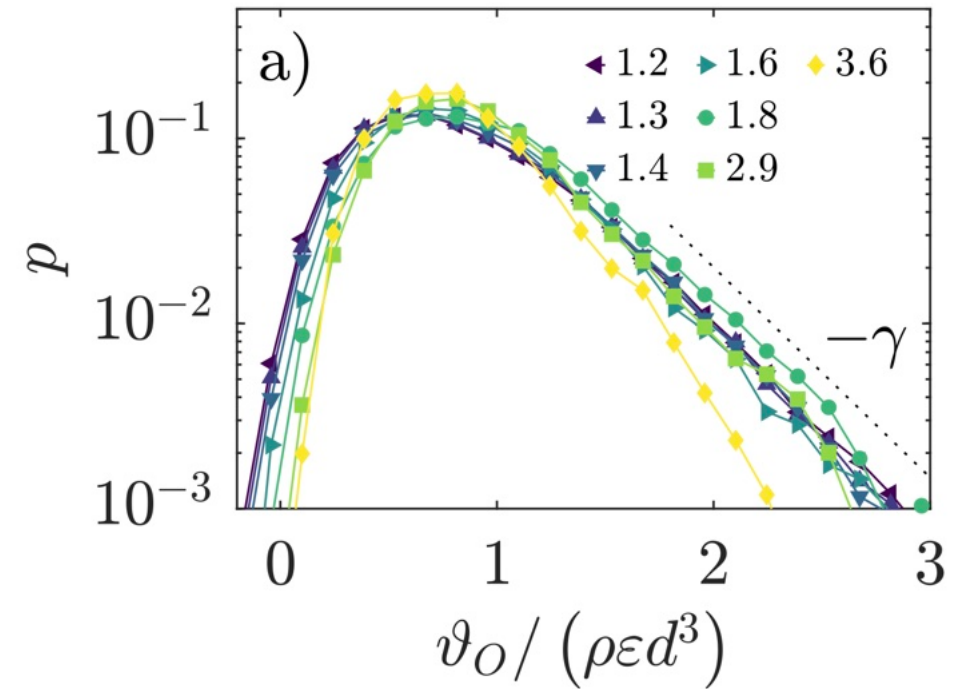
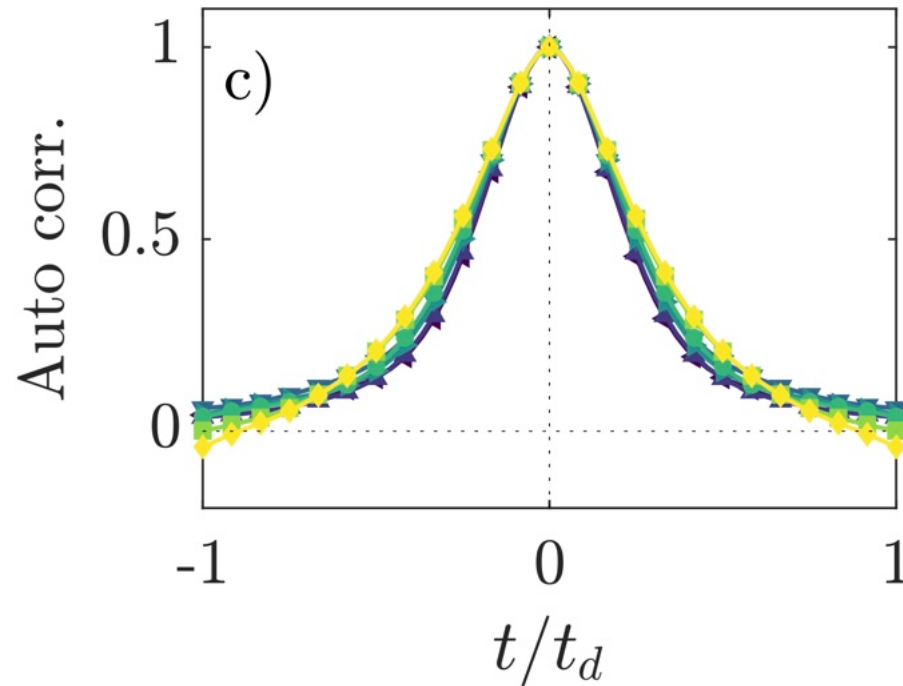
Outer eddies drive deformation

$We = 1.8$



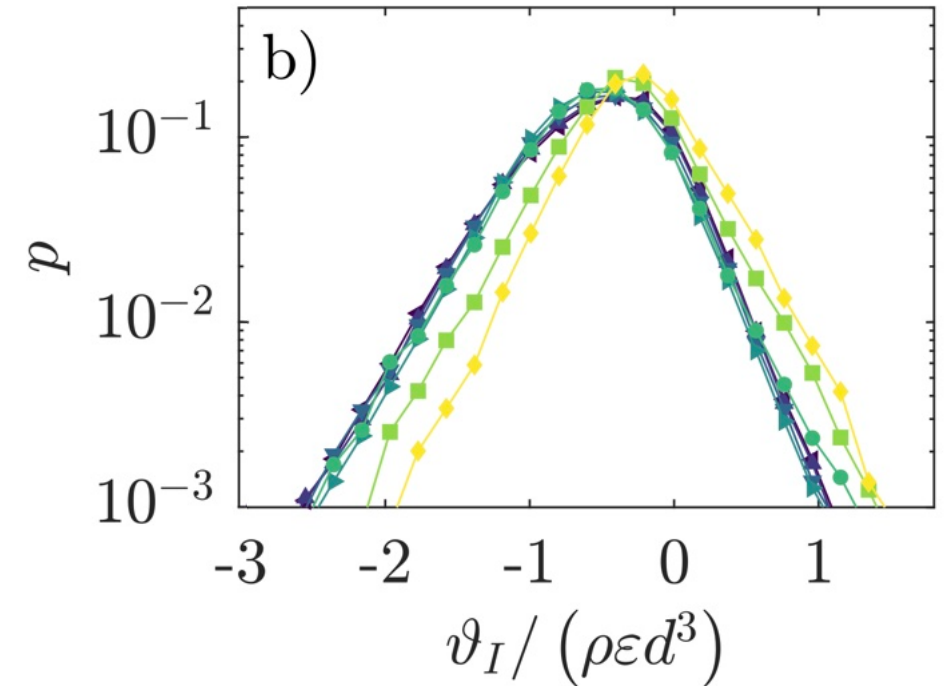
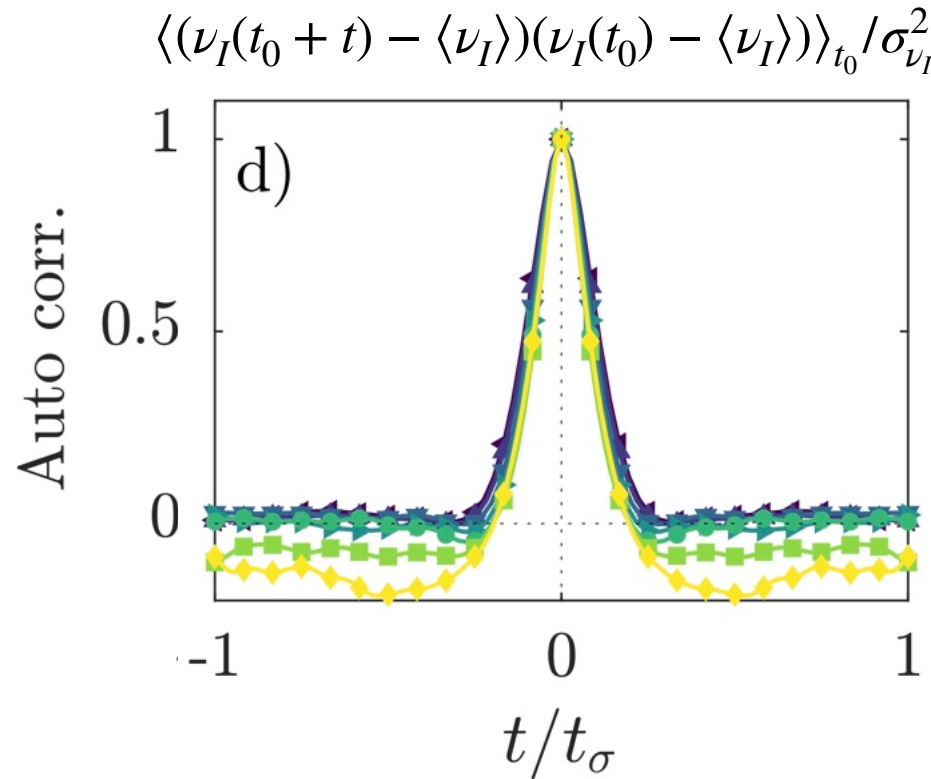
Model of outer stretching

$$\langle (\nu_O(t_0 + t) - \langle \nu_O \rangle)(\nu_O(t_0) - \langle \nu_O \rangle) \rangle_{t_0} / \sigma_{\nu_O}^2$$



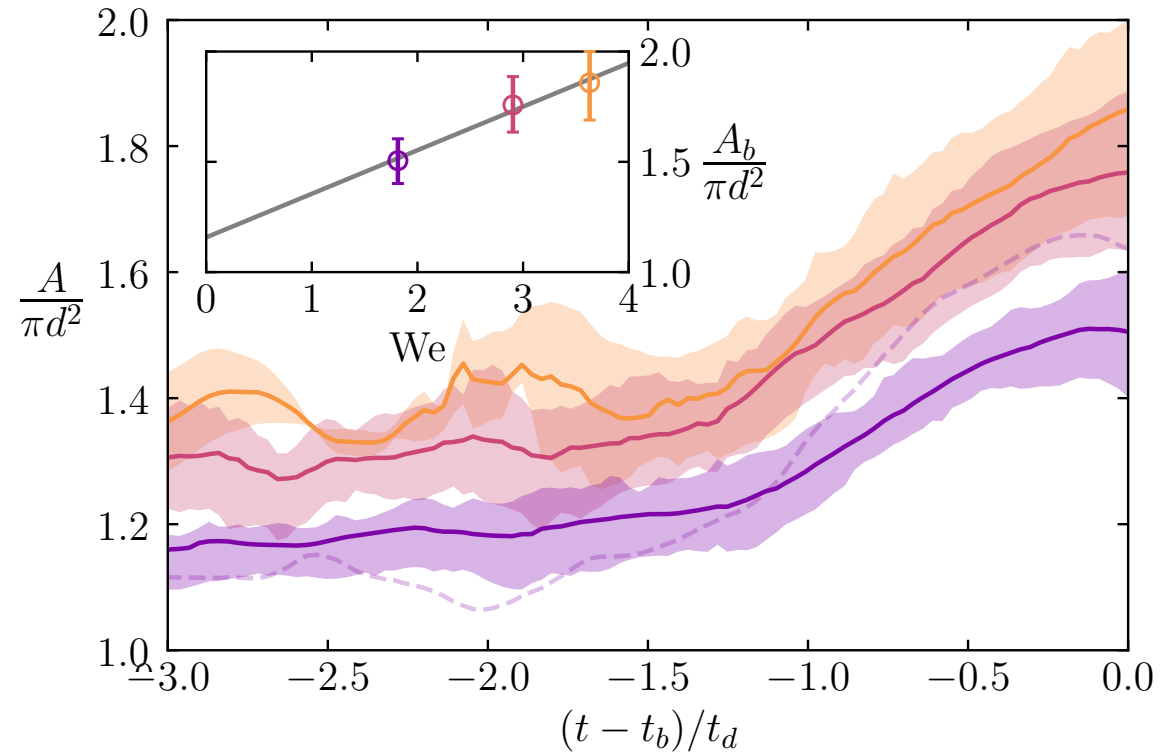
- varies with 0.1 eddy turnover times $t_d = d/u_d$
- value is positive and exponentially distributed. $P(\hat{\vartheta}_O \geq \hat{\vartheta}_O^e) \propto e^{-\gamma \hat{\vartheta}_O^e}, \gamma = 2.6$

Model of inner stretching



- varies with $1/3.5$ capillary times $t_\sigma = \sqrt{\rho d^3 / \sigma}$
- value is negative and depends on the drop area $v_I \approx -\frac{c_d}{t_\sigma} (\mathcal{H} - \mathcal{H}_0)$, $c_d = 3.5$

Area at breakup



Area increase at breakup is constant $\frac{A_b - \langle A \rangle}{\pi d^2} = 0.27$

Breakup rate

- drop energy is a competition $d_t \mathcal{H} = \vartheta_O + \vartheta_I$
- probability of exceeding the breakup energy $P_b \propto e^{-\chi} We^{-3/2}$

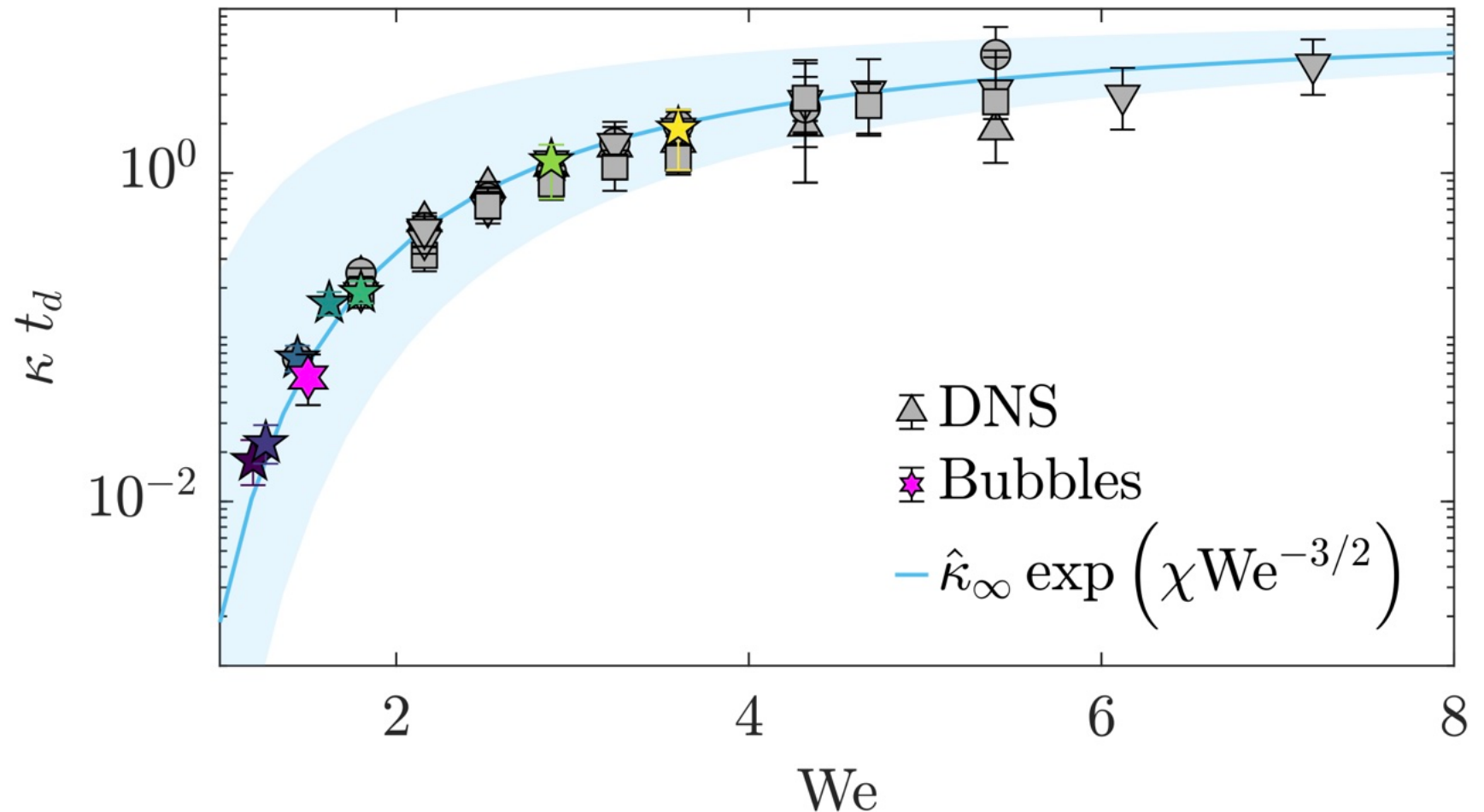
$$\chi = \gamma \pi \xi_c c_d / (1 - \delta)$$

$$\delta = e^{-c_d \hat{t}_{\vartheta_O}} We^{-1/2}$$

- outer stretching events occur ~ 10 times per eddy turnover $\hat{\kappa} = \hat{\kappa}_\infty e^{-\chi} We^{-3/2}$

$$\hat{\kappa}_\infty = 10/t_d$$

Breakup rate



Conclusions



- The drop is stretched by outer eddies and relaxes via inner eddies
- Outer stretching events occur every $0.1t_d$
- Outer stretching intensity is exponentially distributed with $\gamma = 2.6$
- Inner stretching has a rate of $3.5t_\sigma$
- The area increase at breakup is $0.27\pi d^2$
- The breakup rate is $\kappa_\infty t_d^{-1} \exp(-\chi We^{-3/2})$
- Scaling constants can be determined a priori from measurements of the flow
- Model is generalisable to other flows; e.g., non-Newtonian
- Morón, Cannon, Vela-Martín, & Avila, “Causal mechanisms of drop breakup in turbulent flows,” <https://arxiv.org/abs/2605.07504>
- Thank you!